# A Note on the Invariant Distribution of a Stochastic Dynamical System

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#### Abstract

This paper demonstrates the invariant distribution of a stochastic dynamical system. We give the invariant distribution and numerical examples. We also present a further discussion on the computation details.

Keywords Invariant distribution · stochastic dynamical system

# 1 The stochastic dynamic system and invariant distribution

The stochastic dynamic system we focus on is

$$\begin{cases} dx = vdt \\ dv = -\nabla Vdt - \gamma(v - Ax)dt + Avdt + \sigma d\omega, \end{cases}$$
 (1.1)

where

$$V = \frac{1}{2}x^T B x.$$

This equation can be written as

$$d \begin{pmatrix} x \\ v \end{pmatrix} = M \begin{pmatrix} x \\ v \end{pmatrix} dt + \begin{pmatrix} 0 & 0 \\ 0 & I \end{pmatrix} \sigma d\omega,$$

where

$$M = \begin{pmatrix} 0 & I \\ \gamma A - B & A - \gamma I \end{pmatrix}.$$

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We can obtain the solution of the initial form of the problem is

$$x(t) = e^{Mt}x(0) + \int_0^t e^{M(t-\tau)} \begin{pmatrix} 0 & 0 \\ 0 & I \end{pmatrix} d\omega_{\tau}.$$

Remark. We assume that

$$\lim_{t \to +\infty} e^{Mt} = 0.$$

Then x(t) will converge to a limit distribution when  $t \to \infty$ . It can be proved that it is a normal distribution  $\mathcal{N}(0,C)$ , where

$$C = \sigma^2 \int_0^{+\infty} \exp(Mt) \begin{pmatrix} 0 & 0 \\ 0 & I \end{pmatrix} \exp(M^T t) dt.$$

Therefore, we obtain a linear system below

$$MC + CM^{T} = \sigma^{2} \int_{0}^{+\infty} \left( \operatorname{Mexp}(Mt) \begin{pmatrix} 0 & 0 \\ 0 & I \end{pmatrix} \exp\left(M^{T}t\right) + \exp(Mt) \begin{pmatrix} 0 & 0 \\ 0 & I \end{pmatrix} \exp\left(M^{T}t\right) M^{T} \right) dt$$

$$= \sigma^{2} \int_{0}^{+\infty} \frac{d}{dt} \left( \exp(Mt) \begin{pmatrix} 0 & 0 \\ 0 & I \end{pmatrix} \exp\left(M^{T}t\right) \right) dt$$

$$= -\sigma^{2} \begin{pmatrix} 0 & 0 \\ 0 & I \end{pmatrix}.$$

Then we can obtain the theorem below

Theorem 1.1. Suppose that we have an equation

$$\begin{cases} dx = vdt \\ dv = -\nabla Vdt - \gamma(v - Ax)dt + Avdt + \sigma d\omega. \end{cases}$$
 (1.2)

Let

$$M = \left(\begin{array}{cc} 0 & I\\ \gamma A - B & A - \gamma I \end{array}\right)$$

Suppose that

$$\lim_{t \to +\infty} e^{Mt} = 0,$$

then the limit distribution is  $\mathcal{N}(0,C)$  where

$$MC + CM^T = -\sigma^2 \left( \begin{array}{cc} 0 & 0 \\ 0 & I \end{array} \right).$$

From this equation we can not only solve the matrix C easily but also obtain some corollaries. Here we give the most obvious one.

Corollary 1.1. Suppose that

$$C = \left( \begin{array}{cc} C_1 & C_2 \\ C_3 & C_4 \end{array} \right).$$

We have

- (1)  $C_2$  is a skew-symmetric matrix.
- (2) the diagonal elements of  $C_2$  are all 0.

Proof. From

$$MC + CM^T = -\sigma^2 \begin{pmatrix} 0 & 0 \\ 0 & I \end{pmatrix}$$

We can obtain

$$C_2 + C_3 = 0.$$

But C is the covariance matrix of a normal distribution, so C is symmetric. Therefore,

$$C_3 = C_2^T$$

We conclude

$$C_2^T = -C_2$$

So (1) is proved. (2) is a corollary of (1).

### 2 Numerical examples

We give two examples.

Example 2.1.

We obtain

$$c_{11} = \frac{1}{4} \left( 3b^2 + 2 \right), c_{12} = \frac{b}{2}, c_{22} = \frac{1}{2}, c_{33} = \frac{1}{2} \left( b^2 + 1 \right), c_{34} = \frac{b}{4}, c_{44} = \frac{1}{2}, c_{14} = -\frac{b}{4}, c_{44} = \frac{1}{2}, c_{44} = \frac{1}{2}, c_{44} = -\frac{b}{4}, c_{44} = \frac{1}{2}, c_{44} = \frac{$$

Therefore

$$C = \begin{pmatrix} \frac{1}{4} \left(3b^2 + 2\right) & \frac{b}{2} & 0 & -\frac{b}{4} \\ \frac{b}{2} & \frac{1}{2} & \frac{b}{4} & 0 \\ 0 & \frac{b}{4} & \frac{1}{2} \left(b^2 + 1\right) & \frac{b}{4} \\ -\frac{b}{4} & 0 & \frac{b}{4} & \frac{1}{2} \end{pmatrix}.$$

$$\rho = \frac{1}{\sqrt{|2\pi C|}} \exp\left(-\frac{1}{2} \left(q_1, q_2, p_1, p_2\right) C^{-1} \left(q_1, q_2, p_1, p_2\right)^T\right).$$

From the Fokker-Planck equation [1]

$$0 = \frac{d\rho}{dt} = \sum_{i=1}^{n} -\frac{\partial}{\partial q_i} \left( p_i \rho \right) - \frac{\partial}{\partial p_i} \left[ \left( -\frac{\partial V}{\partial q_i} - \gamma \left( p_i - \sum_j A_{ij} q_j \right) + \sum_j A_{ij} p_j \right) \rho \right] + \frac{\sigma^2}{2} \frac{\partial^2}{\partial p_i^2} \rho,$$

we can confirm this result.

#### Example 2.2.

$$n = 1.A = a, B = b$$

$$M = \begin{pmatrix} 0 & 1 \\ \gamma a - b & a - \gamma \end{pmatrix}$$

From

$$MC + CM^T = \left( \begin{array}{cc} 0 & 0 \\ 0 & -\sigma^2 \end{array} \right),$$

we can obtain

$$\begin{split} C = \begin{pmatrix} \frac{\sigma^2}{2(a-\gamma)(\gamma a-b)} & 0\\ 0 & -\frac{\sigma^2}{2(a-\gamma)} \end{pmatrix}\\ \rho = \frac{1}{\sqrt{|2\pi C|}} \exp\left(-\frac{1}{2}(q,p)C^{-1}(q,p)^T\right). \end{split}$$

From the Fokker-Planck equation

$$0 = \frac{d\rho}{dt} = \sum_{i=1}^{n} -\frac{\partial}{\partial q_i} \left( p_i \rho \right) - \frac{\partial}{\partial p_i} \left[ \left( -\frac{\partial V}{\partial q_i} - \gamma \left( p_i - \sum_j A_{ij} q_j \right) + \sum_j A_{ij} p_j \right) \rho \right] + \frac{\sigma^2}{2} \frac{\partial^2}{\partial p_i^2} \rho,$$

## 3 The computation analysis

In this section, we take more discussion on

$$MC + CM^T = -\sigma^2 \left( \begin{array}{cc} 0 & 0 \\ 0 & I \end{array} \right).$$

Assume that  $C' = \frac{1}{\sigma^2}C$ , and in this case, we obtain:

$$MC' + C'M^T = - \left( \begin{array}{cc} 0 & 0 \\ 0 & I \end{array} \right).$$

For conveniently writing, we denote C' as C in the following

$$MC + CM^T = -\begin{pmatrix} 0 & 0 \\ 0 & I \end{pmatrix},$$

where M is known, C is unknown based on simple calculation, we have the conclusion that the total account of the unknown variable is  $2n \times 2n$ . The system is equivalent to a linear equation set DX = y, where  $d \in \Re^{4n^2 \times 4n^2}$ ,  $x \in \Re^{4n^2}$ ,  $y \in \Re^{4n^2}$ . Since

$$C = \left( \begin{array}{cc} C_1 & C_2 \\ -C_2 & C_4 \end{array} \right),$$

where  $C_1$  is a symmetric matrix,  $C_2$  is a skew symmetric matrix,  $C_4$  is a symmetric matrix, the amount of the unknown components, in fact, is only

$$\frac{1}{2}n(n+1) \times 2 + \frac{1}{2}n(n-1) = \frac{3}{2}n^2 + \frac{1}{2}n.$$

Then we obtain that

$$\begin{split} M &= \begin{pmatrix} 0 & I \\ \gamma A - B & A - \gamma I \end{pmatrix}, \\ MC + CM^T &= \begin{pmatrix} 0 & I \\ P & Q \end{pmatrix} \begin{pmatrix} C_1 & C_2 \\ -C_2 & C_4 \end{pmatrix} + \begin{pmatrix} C_1 & C_2 \\ -C_2 & C_4 \end{pmatrix} \begin{pmatrix} 0 & P^T \\ I & Q^T \end{pmatrix} \\ &= \begin{pmatrix} -C_2 & C_4 \\ PC_1 - QC_2 & PC_2 + QC_4 \end{pmatrix} + \begin{pmatrix} C_2 & C_2 P^T + C_2 Q^T \\ C_4 & -C_2 P^T + C_4 Q^T \end{pmatrix} \\ &= \begin{pmatrix} 0 & C_4 + C_1 p^T + C_2 Q^T \\ C_4 + PC_1 - QC_2 & PC_2 + QC_4 - C_2 P^T + C_4 Q^T \end{pmatrix} \\ &= -\begin{pmatrix} 0 & 0 \\ 0 & I \end{pmatrix}, \end{split}$$

where  $P = \gamma A - B$ ,  $Q = A - \gamma I$ .

Thus we obtain that

$$\begin{cases}
C_4 + C_1 P^T + C_2 Q^T = 0, \\
PC_2 + QC_4 - C_2 P^T + C_4 Q^T = I.
\end{cases}$$
(3.1)

According to the first line of (3.1), it holds that

$$C_4 = -C_1 P^T - C_2 Q^T.$$

Besides, we know  $C_4$  is symmetric matrix, which means  $C_4^T = C_4$  Thus we obtain:

$$C_{4} = -C_{1}P^{T} - C_{2}Q^{T}$$

$$= -(C_{1}P^{T} - C_{2}Q^{T})^{T}$$

$$= -PC_{1} - QC_{2}^{T}$$

$$= -PC_{1} - QC_{2}.$$
(3.2)

Substitute (3.2) into the left hand side of the second line of (3.1), and then we obtain that

$$PC_2 + Q(-C_1P^T) - C_2P^T + (PC_1 + QC_2)Q^T = -I,$$

and

$$PC_2 - PC_1Q^T - C_2P^T - QC_1P^T = -I.$$

Thus calculating the covariance matrix C which is equivalent to solving (3.1) is finally equivalent to solving the equations

$$\begin{cases}
C_4 + C_1 P^T + C_2 Q^T = 0 \\
P (C_2 - C_1 Q^T) + (C_2 - C_1 Q^T)^T P^T = -I
\end{cases}$$
(3.3)

without the integral.

#### Example 3.1.

$$n=2, \sigma=1, \gamma=1, A=\left(\begin{array}{cc} 0 & b \\ 0 & 0 \end{array}\right), B=I.$$

so

$$P = \left( \begin{array}{cc} -1 & b \\ 0 & -1 \end{array} \right), Q = \left( \begin{array}{cc} -1 & b \\ 0 & 1 \end{array} \right).$$

We denote  $C_1$  by  $C_{ij}^1$ ,  $C_2$  by  $C_{ij}^2$  and  $C_4$  by  $C_{ij}^4$ . According to symmetry or skew-symmetry of  $C_1, C_2, C_4$ , we obtain the linear equations:

$$\begin{split} -2C_{11}^1 + 2bC_{12}^1 - 2C_{11}^2 + 2b\left(C_{21}^1 - bC_{22}^1 + C_{21}^2\right) &= -1, \\ -C_{12}^1 - C_{21}^1 + bC_{22}^1 - C_{12}^2 - C_{21}^2 + b\left(C_{22}^1 + C_{22}^2\right) &= 0, \\ -C_{11}^1 + bC_{12}^1 - C_{11}^2 + bC_{12}^2 + C_{11}^4 &= 0, \\ -C_{21}^1 + bC_{22}^1 - C_{21}^2 + bC_{22}^2 + C_{21}^4 &= 0, \\ -C_{12}^1 - C_{12}^2 + C_{12}^4 &= 0, \\ -C_{12}^1 - C_{22}^2 + C_{22}^4 &= 0, \\ -2C_{22}^1 - 2C_{22}^2 &= -1, \\ C_{12}^4 - C_{21}^4 &= 0, \\ C_{12}^1 - C_{21}^1 &= 0, \end{split}$$

$$C_{12}^2 + C_{21}^2 = 0,$$
 
$$C_{11}^2 = 0,$$
 
$$C_{22}^2 = 0.$$

Solving the equations, we obtain C as

$$C = \begin{pmatrix} \frac{1}{4} \left(2 + 3b^2\right) & \frac{b}{2} & 0 & -\frac{b}{4} \\ \frac{b}{2} & \frac{1}{2} & \frac{b}{4} & 0 \\ 0 & \frac{b}{4} & \frac{1}{2} \left(1 + b^2\right) & \frac{b}{4} \\ -\frac{b}{4} & 0 & \frac{b}{4} & \frac{1}{2} \end{pmatrix}.$$

Fokker-Planck Equation helps us to check the solution

$$\frac{d\rho}{dt} = \sum_{i=1}^{n} -\frac{\partial}{\partial q_i} \left( p_i \rho \right) - \frac{\partial}{\partial p_i} \left[ \left( -\frac{\partial V}{\partial q_i} - \gamma \left( p_i - \sum_j A_{ij} q_j \right) + \sum_j A_{ij} p_j \right) \rho \right] + \frac{\sigma^2}{2} \frac{\partial^2}{\partial p_i^2} \rho = 0.$$

### 4 Conclusions

This paper demonstrates the invariant distribution of a stochastic dynamical system. We give the invariant distribution and numerical examples. We also give the details of the computation analysis.

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#### References

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